

High-Temperature XRD and DTA Studies of BiMnO₃ Perovskite

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BiMnO₃ synthesized at high pressure is a triclinically distorted perovskite and a ferromagnetic with $T_c = 108$ K. High-temperature X-ray diffraction analysis, thermal analysis, thermogravimetry, and magnetic measurements were used to determine stability and the melting sequence of BiMnO₃ phase in air. Structural phase transition from triclinic to tetragonal structure has been determined at 490°C in air for a polycrystalline BiMnO₃ perovskite. Two phases form immediately at the decomposition temperature of 600°C, namely Bi₂O₃ and Bi₂O₃ · 2Mn₂O_{3+δ}. The Bi₂O₃ · 2Mn₂O_{3+δ} coexists with liquid up to at least 900°C. © 1999 Academic Press

Key Words: perovskite; manganese oxide; high-pressure synthesis; high-temperature XRD; thermal analysis; magnetic property.

1. INTRODUCTION

Among a number of 3d transition metal oxides, the perovskite type manganese oxides, (A_{1-x}B_x)MnO₃, have been studied extensively (1–4). Most studies of manganese oxides with the perovskite structure have concentrated on the properties of the (A_{1-x}B_x)MnO₃ systems containing a larger rare-earth atom such as A = La, Pr, or Nd and a larger alkaline-earth atom such as B = Ca, Sr, or Ba.

Many works led to preparing crystals—varying in quality—of the (A, B)MnO₃ phase using various crystallization methods. Mostly, the crystals are grown from the melt by the flux technique for (La, Pb)MnO₃ (1); however in the cases of A = Pr, B = Ca (3) or A = Sr, Sm, B = Sr (5), the floating zone method is used.

Growing a single crystal of BiMnO₃ is difficult owing to the incongruent character of its melting and also the

very narrow stable temperature range of polycrystalline BiMnO₃. We have tried to develop BiMnO₃ crystals, according to Bokov *et al.* (6), by slow cooling of a completely or partially melted mixture of Bi₂O₃ and Mn₂O₃ in air under normal pressure, but this process failed to produce a BiMnO₃ crystal because of the formation of Bi₂O₃ · 2Mn₂O₃ and Bi₂O₃.

A number of recent works have focussed on the structure and the magnetic properties of BiMnO₃ (4, 7). The results available up to now are:

- (i) BiMnO₃ is ferromagnetic with a Curie temperature of 105–110 K (4, 6, 8);
- (ii) BiMnO₃ can be indexed on the basis of a triclinic unit cell;
- (iii) A polycrystalline BiMnO₃ prepared under high pressure decomposes when it is heated at ordinary pressure. The upper limit of this decomposition temperature is 300°C in vacuum (8). On the other hand, Tomashpol'skii *et al.* reported that polycrystalline BiMnO₃ shows a transition at 210°C in relation to the change of the superstructure and then begins to decompose at 600°C (9).

The results given here correspond to the description of the phase equilibrium by investigation in detail of the effect of temperature. In order to produce a BiMnO₃ single crystal under normal pressure, more information about the stability of BiMnO₃ is desirable. Ideally, high-temperature X-ray diffraction (HTXRD) gives phase information *in situ*, without the difficulties encountered in quench studies. However, there is poor agreement among the HTXRD results in the literature, especially for the BiMnO₃ phase. Therefore, we have undertaken a detailed stability study of BiMnO₃ perovskite using HTXRD, thermal analysis, and magnetic measurement. In this article, we report a comprehensive study on the decomposition and stability of BiMnO₃ perovskite in air.

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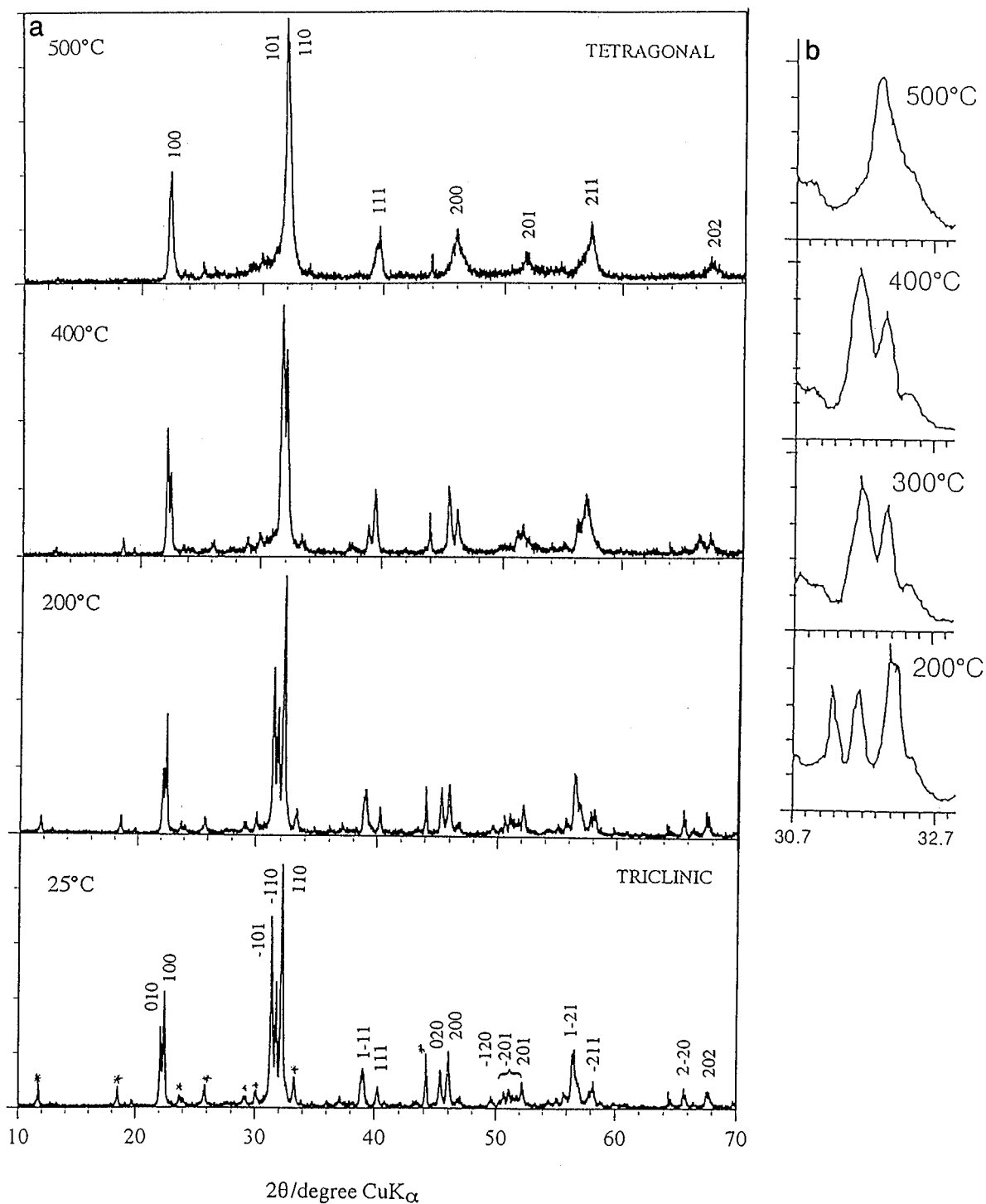


FIG. 1. High temperature X-ray diffraction for the BiMnO₃ perovskite. (a) From 25 to 500°C (the asterisk indicates peaks of impurity compounds). (b) Part of the high temperature X-ray diffraction showing the structural change of BiMnO₃ perovskite from triclinic to tetragonal near 500°C. (c) From 600 to 900°C.

2. EXPERIMENTAL CONDITIONS

The synthesis of a polycrystalline BiMnO₃ sample was carried out by using a high-pressure technique. The starting

materials used were Bi₂O₃ and Mn₂O₃; a mixture of the necessary stoichiometry was fired at 300°C for 24 h before use. The mixture was sealed in a gold capsule (∅ 2.5 mm) to avoid interaction with surroundings and placed in a

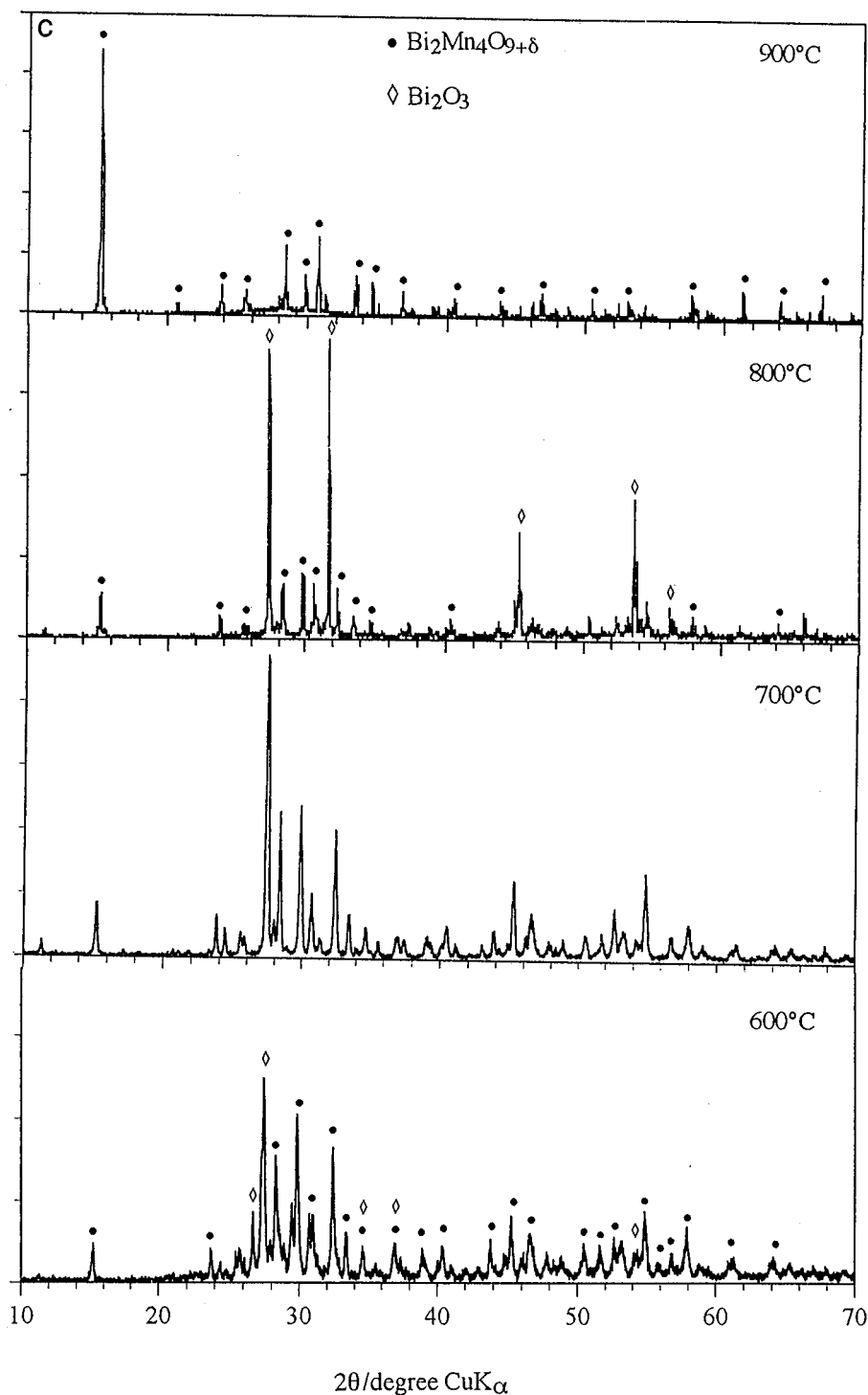


FIG. 1—Continued

cylindrical graphite heater electrically insulated by a BN sleeve and then was placed in an 8 mm pyrophyllite cube container which served as the pressure transmission medium.

The BiMnO_3 perovskite was synthesized under 6 GPa in the temperature range from 600 to 610°C and was kept for

3 h and then quenched (4). Temperature was measured by a Pt/Pt-13% Rh thermocouple. The pressure was applied to the sample first and then the temperature was brought to the desired value and held for a predetermined interval of time. The pressure was removed gradually after quenching.

The prepared sample was examined by HTXRD in a Philips 1700 diffractometer, using $\text{CuK}\alpha$ radiation at 45 kV and 30 mA.

To determine the range of stability of the BiMnO_3 phase, differential thermal analysis and thermogravimetry (TG/DTA 92-Setaram) was performed in an Al_2O_3 crucible using a reference material. The temperature was measured by a thermocouple placed at the bottom of each crucible.

The magnetization (M) was measured in the temperature range 5–350 K with a superconducting quantum interference device magnetometer (Quantum Design). The susceptibility measurements were performed in a 1 T field.

3. RESULTS AND DISCUSSION

3.1. X-ray Diffraction

The room temperature X-ray patterns of the BiMnO_3 perovskite are shown in Fig. 1a. BiMnO_3 crystallize in the triclinic structure at room temperature, all reflections could be indexed on the basis of the triclinic cell. The lattice parameters refined by least-squares analysis are $a = 3.950 \text{ \AA}$, $b = 3.995 \text{ \AA}$, $c = 3.919 \text{ \AA}$, $\alpha = 90.7^\circ$, $\beta = 90.9^\circ$, and $\gamma = 91.9^\circ$.

The indexed powder diffraction data are listed in Table 1 and the X-ray powder diffraction patterns comparing low (at 25°C) and high forms (at 500°C) are illustrated in Fig. 1. These data were collected using a $0.3^\circ/\text{min}$ scanning rate. A small amount of bismuth oxycarbonate, $\text{Bi}_2(\text{CO}_3)\text{O}_2$, was found in the powder XRD pattern at room temperature. The carbon atoms are considered to come from the graphite heater through narrow cracks in the gold capsule during the heating process.

Structural transformation of BiMnO_3 with increasing temperature from room temperature to 500°C was observed. The unit cell dimensions determined from the observed d -spacings by least-squares are summarized in Table 1 together with the firing conditions. In order to confirm the thermally induced structural phase transition from triclinic to tetragonal, the change of the (-110) , (-101) , and (110) peak intensity between $2\theta = 30.7$ and 31.7 (where θ is diffraction angle) was measured out with a step of 0.02° and a scanning rate of $0.001^\circ/\text{s}$. The change of the triclinic (hkl) peak intensity with temperature was traced in Fig. 1b.

At 600°C , the corresponding peaks of the BiMnO_3 phase disappeared, indicating that BiMnO_3 polycrystalline is unstable in air for temperatures above 600°C (Fig. 1c). At the same temperature Bi_2O_3 and $\text{Bi}_2\text{O}_3 \cdot 2\text{Mn}_2\text{O}_{3+\delta}$ were the most prominent phases. The intensities of diffraction peaks corresponding to the Bi_2O_3 and $\text{Bi}_2\text{O}_3 \cdot 2\text{Mn}_2\text{O}_{3+\delta}$ phases become stronger as temperature increased. The phase content versus temperature is shown in Fig. 2, which gives the

TABLE 1
Observed and Calculated d -Spacings and Intensities of BiMnO_3 :
(a) at 25°C Low Form^a and (b) at 500°C High Form^b

h	k	l	I_{obs}	I_{cal}	$d_{\text{obs}} (\text{\AA})$	$d_{\text{cal}} (\text{\AA})$
a						
0	1	0	31	28	4.004	3.994
1	0	0	46	29	3.967	3.958
-1	0	1	74	93	2.818	2.813
-1	1	0	47	47	2.856	2.844
1	1	0	100	100	2.779	2.783
1	-1	1	25	35	2.303	2.292
1	1	1	11	32	2.243	2.247
0	2	0	15	50	2.001	1.997
2	0	0	25	52	1.974	1.974
1	-2	1	21	17	1.634	1.634
-2	1	1	24	18	1.626	1.628
b						
1	0	0	42	23	3.949	3.943
1	0	1	92	45	2.796	2.784
1	1	0	100	100	2.789	2.788
1	1	1	17	26	2.268	2.274
2	0	0	17	26	1.973	1.972
2	0	1	8	14	1.768	1.762
2	1	1	19	9	1.605	1.609
2	0	2	5	7	1.392	1.392

^a Calculated on the basis of a triclinic unit cell: $a = 3.950 \text{ \AA}$, $b = 3.995 \text{ \AA}$, $c = 3.919 \text{ \AA}$, $\alpha = 90.7^\circ$, $\beta = 90.9^\circ$, and $\gamma = 91.0^\circ$.

^b Calculated on the basis of a tetragonal unit cell: $a = b = 3.943 \text{ \AA}$, $c = 3.930 \text{ \AA}$, $\alpha = \beta = \gamma = 90.0^\circ$.

peak height intensities for each phase in the melting sequence shown in Fig. 1. Note that the peak intensities do not give quantitative amounts of each phase in the sample.

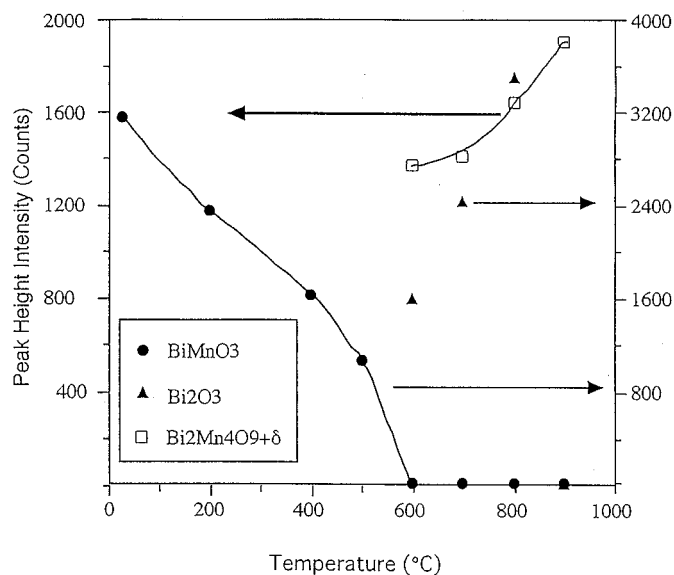
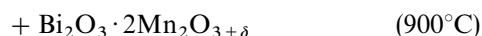
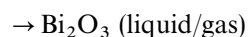
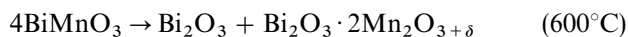


FIG. 2. Relative peak intensities of BiMnO_3 , Bi_2O_3 , and $\text{Bi}_2\text{O}_3 \cdot 2\text{Mn}_2\text{O}_{3+\delta}$ compounds versus temperature between 25 and 900°C .

So the composition cannot be deduced from the data in Fig. 2. However, the relative peak intensities of the same phase give a reasonable approximation of the amount of that phase versus temperature. The relative peak intensities corresponding to BiMnO₃ decrease with temperature between 25 and 500°C; this phenomenon is caused by thermal vibration.

At $T = 600^\circ\text{C}$ the Bi₂O₃ phase appears and grows faster than Bi₂O₃·2Mn₂O_{3+δ}. The Bi₂O₃ phase shows a structural change to cubic phase at $T = 800^\circ\text{C}$. All peaks could be indexed on the basis of the cubic lattice with $a = 5.66 \text{ \AA}$ according to (10). The evolution sequence of Bi₂O₃ determined by HTXRD agrees qualitatively with the phase diagram reported for the Bi₂O₃-Mn₂O₃ system (11). The melting starts on reaching 830°C, but Bi₂O₃·2Mn₂O_{3+δ} phase coexists with a liquid up to at least 900°C (Fig. 1c), which was the maximum temperature investigated. At 900°C the diffraction peaks corresponding to the Bi₂O₃ vanished and the major peaks were identified to be

Bi₂O₃·2Mn₂O_{3+δ} (12). High temperature X-ray diffraction suggested the following process:



3.2. Thermal Analysis

The measurement of weight change and differential thermal analysis were done by heating a 37 mg specimen to 950°C and then cooling it to room temperature at a constant rate of 5°C/min. The TG/DTA traces are shown in Fig. 3. DTA measurements showed three peaks at 490, 630, and 740°C. The endothermic DTA peak observed near 490°C is likely to correspond to phase transition BiMnO₃

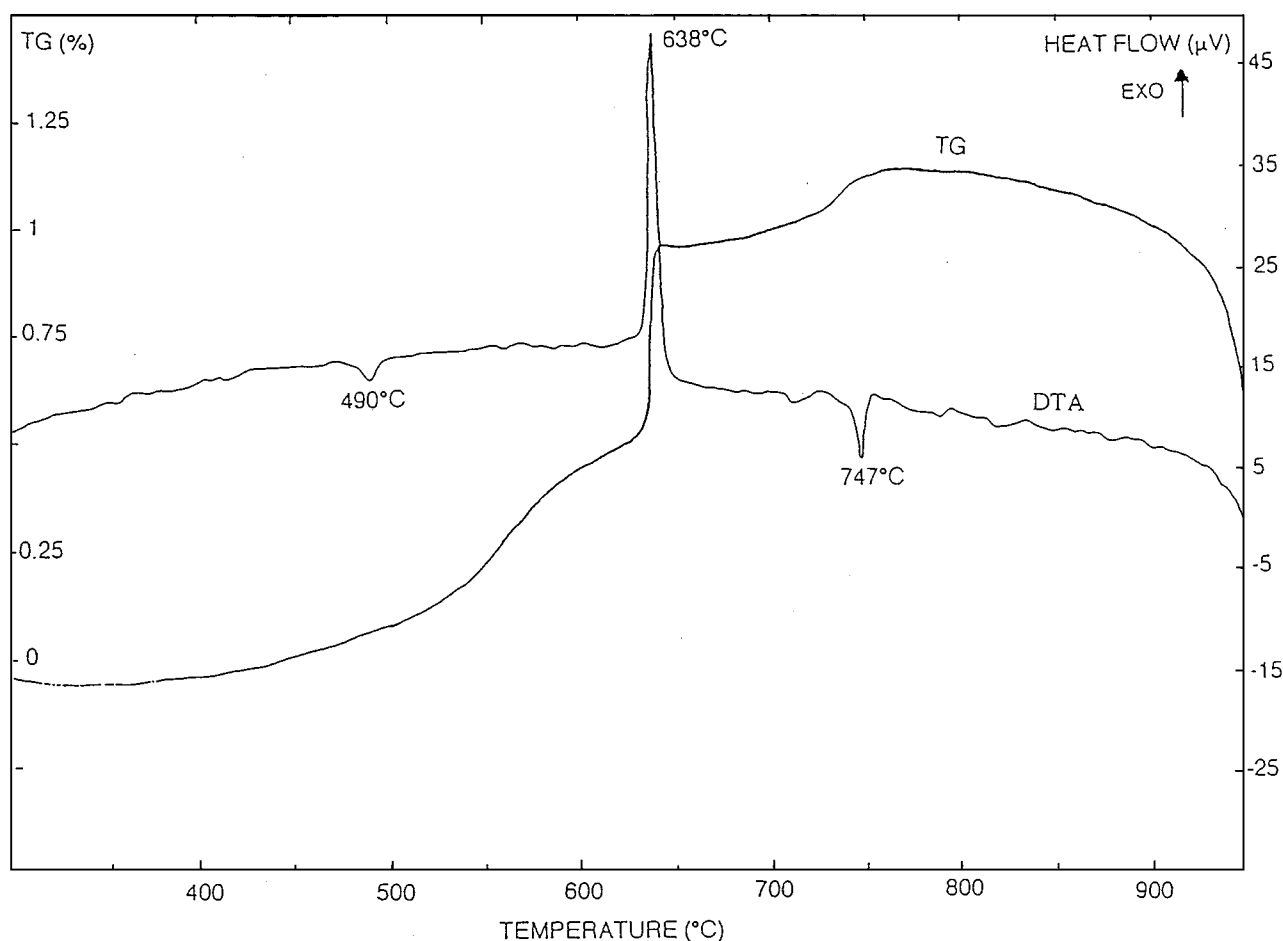


FIG. 3. TG/DTA curves showing thermal behavior of a BiMnO₃ sample from 350 to 950°C.

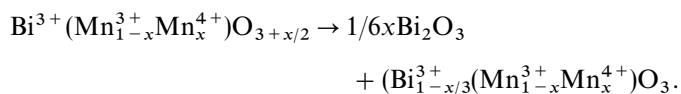
from triclinic to tetragonal, whereas the second one, at 740°C, can be attributed to phase transition Bi_2O_3 from monoclinic to cubic according to the pseudobinary $\text{Bi}_2\text{O}_3 \cdot \text{Mn}_2\text{O}_3$ system (11).

The lower stability limit was detectable by DTA measurements, which showed a clear exothermal reaction at 630°C combined with a weight increase of around 0.6% at this temperature, whose stability seems to depend severely on the oxygen partial pressure. The mass of sample increased during heating in three steps, at 480–610°C, 620–640°C, and 640–770°C. To explore the first transition at 490°C, the sample of BiMnO_3 was heated to 500°C in the TG/DTA measurements and maintained at this temperature for 30 min and then cooled to room temperature at a constant rate of 5°C/min. DTA revealed no exothermal peak during cooling, indicating that the structural change of BiMnO_3 from triclinic to tetragonal is not reversible. Moreover, the X-ray powder diffraction patterns of the compound at room temperature and the quenched sample from 500°C in air are different.

There are two possible explanations for weight change between room temperature and 500°C: the oxidation state of the bismuth has partly increased from 3+ to 5+ ($\text{Bi}_{1-x}^{3+}\text{Bi}_x^{5+}\text{Mn}^{3+}\text{O}_{3+x}$), which would decrease the ionic radius of bismuth, from 0.96 Å (Bi^{3+}) to 0.74 Å (Bi^{5+}), explaining the decrease in volume (see Table 1), or the oxidation state of Mn has increased from 3+ to 4+ ($\text{Bi}^{3+}\text{Mn}_{1-x}^{3+}\text{Mn}_x^{4+}\text{O}_{3+x/2}$). Since Mn^{3+} is a Teller ion and Mn^{4+} is not, the change in oxidation would remove the Teller distortion, and this would also explain the decrease in ionic radii from Mn^{3+} to Mn^{4+} . Magnetic measurements of BiMnO_3 before and after annealing at 500°C should allow us to determine the oxidation state of Mn and determine which of these possibilities is correct.

Assuming the existence of such nonstoichiometric perovskite, at least two schemes of behavior are possible:

(a) Interstitial oxygen may be present, occupying probably the (1/2 0 0) sites of low electrostatic potential (of the primitive perovskite cell) or perhaps the smaller tetrahedral site (1/4 1/4 1/4) (13, 14). In this case we may write



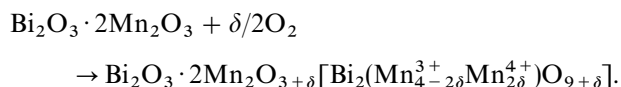
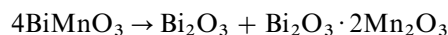
In such a situation, however, we would expect to detect the Bi_2O_3 phase by X-ray diffraction analysis. A region of oxidative nonstoichiometry, $0 < x < 0.15$, for LaMnO_3 was reported by Wold *et al.* (15) with a change in symmetry from orthorhombic to rhombohedral for Mn^{4+} concentrations greater than 21% ($\text{LaMnO}_{3.10}$). On the other hand, Verselst *et al.* (16) reported that the structure of $\text{LaMnO}_{3+\delta}$ changed from orthorhombic to cubic, via rhombohedral, with increased Mn^{4+} content (20% Mn^{4+}). The structural phase transition between the orthorhombic and rhombohedral

was observed neither in a single crystal of perovskite-type manganese oxide ($\text{La}_{1-x}\text{Sr}_x\text{MnO}_3$ ($x = 0.17$) (2) nor in ($\text{Bi}_{1-x}\text{Sr}_x\text{MnO}_3$), which reveals a triclinic-to-tetragonal change (via monoclinic) with partial replacement of bismuth by strontium ($x = 0.5$) (4).

(b) Alternatively, vacancies may exist at both Bi and Mn cation sites. Recently the cation vacancies of $\text{La}_{1-\delta}\text{Mn}_{1-\delta}\text{O}_3$ have been reported by many groups (14, 17, 18).

To discriminate between these various models, only structural resolution of a BiMnO_3 single crystal can indicate to us the exact composition. In order to solve this problem, we are going to obtain a BiMnO_3 single crystal and then make a structural refinement.

The uptake in oxygen from 630 to 770°C increases the weight by about 0.7%. The weight increase is probably due to partial oxidation of Mn^{3+} to Mn^{4+} , assuming the following reactions:



The average manganese valency is estimated, from thermogravimetric analysis, to be 3.15 for $\text{Bi}_2\text{O}_3 \cdot 2\text{Mn}_2\text{O}_{3+\delta}$. This corresponds to 18% of the manganese atoms in the 4+ state. The weight loss observed at 900°C can be attributed to Bi_2O_3 evaporation. The temperature of the thermal events are consistent with HTXRD data.

3.3. Magnetic Measurement

The temperature dependence of magnetization and inverse of molar magnetic susceptibility ($1/\chi_m$) of BiMnO_3 are shown in Fig. 4. BiMnO_3 synthesized at high pressure is a ferromagnet with T_c about 108 K (Fig. 4a), consistent with the previous work (4, 8). The ferromagnetism of BiMnO_3 has been explained by analyzing the Mn–O–Mn angles (19). To clarify the origin of the ferromagnetism in the BiMnO_3 sample, accurate structure resolution of a single crystal is indispensable. The magnetization measured in a field of 1 T is 1.72 emu/Oe·mol, giving a magnetic moment of 3.1 μ_B per Mn atom at 5 K, comparable with the full spin value of Mn^{3+} . The relationship of $1/\chi_m$ vs T obeys the Curie–Weiss law above T_c . The effective Bohr magneton number estimated from the measured slope of $1/\chi_m$ vs T also gives a value of $p_{\text{eff}} = 5.16$, which is in very good agreement with the spin-only value of Mn^{3+} ($p_{\text{so}} = 4.90$).

Annealing of ferromagnetic BiMnO_3 at 500°C for 4 h drastically reduces its ferromagnetic moment, as shown in Fig. 4b. A very weak magnetic moment appears below about 80 K with paramagnetic Curie temperature (θ_p) \approx 50 K. The effective magnetic moments p_{eff} extracted from

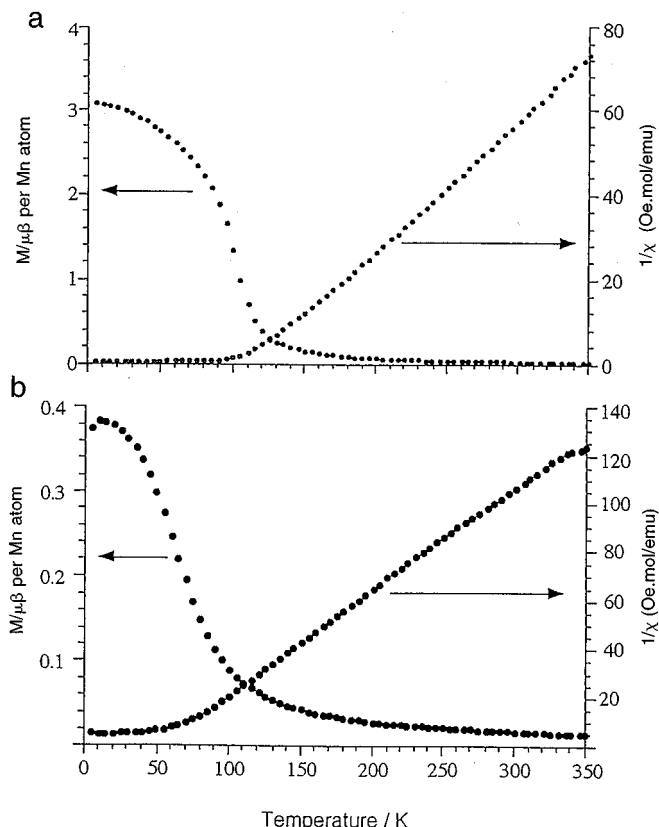


FIG. 4. Temperature variation of magnetization and inverse of molar magnetic susceptibility of BiMnO₃ perovskite (measured at 1 T) (a) before annealing and (b) after annealing at 500°C for 4 h.

the inverse susceptibility decreased from 5.16 to 4.37, suggesting that the oxidation state of some of the manganese has increased from 3+ to 4+ after annealing the sample at 500°C.

The observed remarkable change in magnetic properties of BiMnO₃ with annealing corresponds to the structure change from triclinic to tetragonal, as is observed in the (Bi, Sr)MnO₃ system (4).

4. CONCLUSION

We have investigated a BiMnO₃ perovskite, prepared under high pressure (6 GPa), using high-temperature X-ray diffraction analysis, thermal analysis, and magnetic measurements. We can propose a chemical path for the

thermal behavior of BiMnO₃ perovskite from room temperature to 900°C.

(i) The structure of BiMnO₃ compound changes from triclinic to tetragonal near 490°C. This structural change was found to be irreversible;

(ii) BiMnO₃ undergoes peritectoidic decomposition around 630°C with formation of Bi₂O₃ and Bi₂O₃ · 2Mn₂O₃;

(iii) There is a weight gain around 630°C, with a partial oxidation of Mn³⁺ to Mn⁴⁺;

(iv) The Bi₂O₃ phase disappears during heating and the Bi₂O₃ · 2Mn₂O_{3+δ} phase coexists with a liquid up to 900°C.

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REFERENCES

1. A. H. Morrish, B. J. Evans, J. A. Eaton, and L. K. Leung, *Can. J. Phys.* **47**, 2691 (1969).
2. A. Asamitsu, Y. Moritomo, R. Kumai, Y. Tomioka, and Y. Tokura, *Phys. Rev. B* **54**, 1716 (1996).
3. Y. Tomioka, Y. Asamitsu, H. Kuwahara, Y. Moritomo, and Y. Tokura, *Phys. Rev. B* **53**, R1689 (1996).
4. H. Chiba, T. Atou, and Y. Syono, *J. Solid State Chem.* **132**, 139 (1997).
5. A. Urushibara, Y. Moritomo, T. Arima, A. Asamitsu, G. Kido, and Y. Tokura, *Phys. Rev. B* **51**, 14103 (1995).
6. V. A. Bokov, I. E. Myl'nikova, S. A. Kizhaev, M. F. Bryzhina, and N. A. Grigoryan, *Sov. Phys. Solid State* **7**, 2993 (1966).
7. I. O. Troyanchuk, N. V. Kasper, O. S. Mantyskaya, and S. N. Pastushonok, *Sov. Phys. JETP* **78**(2), 212 (1994). [English translation]
8. F. Sugawara, S. Iida, Y. Syono, and S. Akimoto, *J. Phys. Soc. Jpn.* **25**, 1553 (1968).
9. Yu. Ya. Tomashpol'skii, E. V. Zubova, K. P. Burdina, and Yu. N. Venetsev, *Izv. Akad. Nauk SSSR, Inorg. Mater.* **3**, 2132 (1967).
10. Gattow Schroder, *Z. Anorg. Allg. Chem.* **318**, 176 (1962).
11. E. M. Levin and R. S. Roth, *J. Res. Natl. Bur. Standards* **68A** (2), 197 (1964).
12. N. Niizeki and M. Wachi, *Z. Kristallogr.* **127**, 173 (1968).
13. B. C. Tofield and W. R. Scott, *J. Solid State Chem.* **10**, 183 (1974).
14. J. A. M. Van Roosmalem, E. H. P. Cordefunke, R. B. Helmholtz, and H. W. Zandbergen, *J. Solid State Chem.* **110**, 100 (1994).
15. A. Wold and R. J. Arnott, *J. Phys. Chem. Solids* **9**, 176 (1959).
16. M. Verselst, N. Rangavittal, C. N. R. Rao and A. Rousset, *J. Solid State Chem.* **104**, 74 (1993).
17. J. Töpfer and J. B. Goodenough, *J. Solid State Chem.* **130**, 117 (1997).
18. A. Arulraj, R. Mahesh, G. N. Subbanna, R. Mahendiran, A. K. Raychaudhuri, and C. N. R. Rao, *J. Solid State Chem.* **127**, 87 (1996).
19. E. E. Havinga, *Philips Res. Rep.* **21**, 432 (1966).